

Orthogonality of mode shapes for an Euler-Bernoulli beam

For the r th normal mode $\frac{d^2}{dx^2} \left[EI(x) \frac{d^2 \phi_r}{dx^2} \right] = \omega_r^2 m(x) \phi_r(x)$ — (1)

For the s th normal mode $\frac{d^2}{dx^2} \left[EI(x) \frac{d^2 \phi_s}{dx^2} \right] = \omega_s^2 m(x) \phi_s(x)$ — (2)

Multiplying both sides of Eqn. (1) with $\phi_s(x)$ and integrating, we get,

$$\int_0^L \phi_s \frac{d^2}{dx^2} \left[EI(x) \frac{d^2 \phi_r}{dx^2} \right] dx = \omega_r^2 \int_0^L \phi_s m(x) \phi_r dx$$

$= 0$ (for free, ss, fixed) _{BCs}

$$\Rightarrow \cancel{\phi_s \frac{d}{dx} \left[EI(x) \frac{d^2 \phi_r}{dx^2} \right]} \Big|_0^L - \int_0^L \frac{d\phi_s}{dx} \frac{d}{dx} \left[EI(x) \frac{d^2 \phi_r}{dx^2} \right] dx = \omega_r^2 \int_0^L \phi_s m(x) \phi_r dx$$

$= 0$ (for free, fixed, ss BCs)

$$\Rightarrow \cancel{-\frac{d\phi_s}{dx} EI(x) \frac{d^2 \phi_r}{dx^2}} \Big|_0^L + \int_0^L \frac{d^2 \phi_s}{dx^2} EI(x) \frac{d^2 \phi_r}{dx^2} dx = \omega_r^2 \int_0^L \phi_s m(x) \phi_r dx$$

$$\Rightarrow \int_0^L \frac{d^2 \phi_s}{dx^2} EI(x) \frac{d^2 \phi_r}{dx^2} dx = \omega_r^2 \int_0^L \phi_s(x) m(x) \phi_r(x) dx$$
 — (3)

Similarly, we can get, (by multiplying both sides of Eqn. 2 with $\phi_r(x)$)

$$\int_0^L \frac{d^2 \phi_r}{dx^2} EI(x) \frac{d^2 \phi_s}{dx^2} dx = \omega_s^2 \int_0^L \phi_r(x) m(x) \phi_s(x) dx$$
 — (4)

Subtracting Eqn. (4) from Eqn. (3), we get,

$$(\omega_r^2 - \omega_s^2) \int_0^L \phi_r(x) m(x) \phi_s(x) dx = 0$$

for $r \neq s$, $\omega_r^2 \neq \omega_s^2 \Rightarrow \int_0^L \phi_r(x) m(x) \phi_s(x) dx = 0$, $\int_0^L \phi_r(x) m(x) \phi_r(x) dx = 1$
 (for mass normalized mode shapes)

and

$$\int_0^L \frac{d^2 \phi_r}{dx^2} EI(x) \frac{d^2 \phi_s}{dx^2} dx = 0 \quad \text{and} \quad \int_0^L \frac{d^2 \phi_r}{dx^2} EI(x) \frac{d^2 \phi_r}{dx^2} dx = \omega_r^2$$

Fixed-free Euler-Bernoulli beam:

$$\phi(x) = c_1 \sin \beta x + c_2 \cos \beta x + c_3 \sinh \beta x + c_4 \cosh \beta x$$

$$\textcircled{a} x=0 \quad y(0,t) = \phi(0)T(t) = 0 \Rightarrow \phi(0) = 0$$

$$\textcircled{a} x=0 \quad y'(0,t) = \phi'(0)T(t) = 0 \Rightarrow \phi'(0) = 0$$

$$\phi(0) = 0 \Rightarrow c_2 + c_4 = 0 \quad \text{and} \quad \phi'(0) = c_1 + c_3 = 0$$

$$\therefore \phi(x) = c_1 (\sin \beta x - \sinh \beta x) + c_2 (\cos \beta x - \cosh \beta x)$$

$$\textcircled{a} x=L, \quad M = EI \frac{\partial^2 y}{\partial x^2} \Rightarrow EI \phi''(L)T(t) = 0 \Rightarrow \phi''(L) = 0$$

$$\textcircled{a} x=L \quad V = EI \frac{\partial^3 y}{\partial x^3} = EI \phi'''(L)T(t) = 0 \Rightarrow \phi'''(L) = 0$$

$$-\beta^2 c_1 (\sin \beta L + \sinh \beta L) - \beta^2 c_2 (\cos \beta L + \cosh \beta L) = 0$$

$$-\beta^3 c_1 (\cos \beta L + \cosh \beta L) + \beta^3 c_2 (\sin \beta L - \sinh \beta L) = 0$$

$$\Rightarrow \begin{bmatrix} \sin \beta L + \sinh \beta L & \cos \beta L + \cosh \beta L \\ \cos \beta L + \cosh \beta L & -(\sin \beta L - \sinh \beta L) \end{bmatrix} \begin{Bmatrix} c_1 \\ c_2 \end{Bmatrix} = 0$$

$$\Rightarrow \boxed{\cos \beta L \cosh \beta L = -1} \Rightarrow \beta L = \begin{matrix} 1.8751 \\ 4.6941 \\ 7.8548 \\ 10.996 \end{matrix}$$

Modal analysis : Solution of continuous system

Eqn. of motion of rod under axial vibration

$$\frac{\partial}{\partial x} \left[EA(x) \frac{\partial u}{\partial x} \right] + f(x,t) = m(x) \frac{\partial^2 u}{\partial t^2} \quad \text{--- (1)}$$

Modal expansion theorem

$$u(x,t) = \sum_{r=1}^{\infty} \phi_r(x) \eta_r(t) \quad \text{--- (2)}$$

Substituting Eqn. (2) into Eqn. (1), we get,

$$\frac{d}{dx} \left[EA(x) \sum_r \phi_r'(x) \eta_r(t) \right] + f(x,t) = m(x) \sum_r \phi_r(x) \ddot{\eta}_r(t) \quad \text{--- (3)}$$

Multiplying both sides of Eqn. (3) with $\phi_s(x)$ and integrating, we get,

$$\int_0^L \phi_s(x) \frac{d}{dx} \left[EA(x) \sum_r \phi_r'(x) \eta_r(t) \right] dx + \int_0^L \phi_s(x) f(x,t) dx$$

$$= \int_0^L \phi_s(x) m(x) \sum_r \phi_r(x) \ddot{\eta}_r(t) dx = \ddot{\eta}_r(t) \text{ for } r=s$$

$$= 0 \text{ for } r \neq s$$

$$\Rightarrow \phi_s(x) EA(x) \sum_r \phi_r'(x) \eta_r(t) \Big|_0^L - \int_0^L \frac{d\phi_s}{dx} EA(x) \frac{d\phi_r}{dx} \eta_r(t) dx = \ddot{\eta}_r(t)$$

= 0 for all BCs
(fixed/free/spring)

$$+ \int_0^L \phi_s(x) f(x,t) dx$$

$$\Rightarrow - \int_0^L \frac{d\phi_s}{dx} EA(x) \frac{d\phi_r}{dx} \eta_r(t) dx + \int_0^L \phi_s(x) f(x,t) dx = \ddot{\eta}_r(t)$$

$\overset{= \omega_r^2 \eta_r(t)}{= 0}$ (for $s=r$)
 $\overset{= 0}{= 0}$ (for $s \neq r$)

$$\Rightarrow -\eta_r(t) \omega_r^2 + \int_0^L \phi_r(x) f(x,t) dx = \ddot{\eta}_r(t)$$

$$\Rightarrow \ddot{\eta}_r(t) + \omega_r^2 \eta_r = \int_0^L \phi_r(x) f(x,t) dx = Q_r(t)$$

Example

Consider a fixed-fixed rod under axial vibration. Derive the expression for the response of the rod due to an initial condition given as $u(x,0) = Ax(1 - \frac{x}{L})$

Step 1

Natural frequencies
&
mode shapes.

$$\phi(x) = C_1 \sin \beta x + C_2 \cos \beta x$$

BCs

$$\text{@ } x=0, \phi(0)=0 \Rightarrow C_2=0$$

$$\text{@ } x=L, \phi(L)=0 \Rightarrow C_1 \sin \beta L = 0 \Rightarrow \beta = \frac{\lambda r}{L} \Rightarrow \omega_r = \frac{\lambda r}{L} \sqrt{\frac{EA}{m}}$$

and $\phi_r(x) = C_r \sin \frac{\lambda r x}{L}$, Mass normalized mode shape. $\phi_r(x) = \sqrt{\frac{2}{mL}} \sin \frac{\lambda r x}{L}$

Step 2

Eqn. of motion after modal analysis

$$\ddot{\eta}_r + \omega_r^2 \eta_r = 0$$

$$\Rightarrow \eta_r(t) = \eta_r(0) \cos \omega_r t + \frac{\dot{\eta}_r(0)}{\omega_r} \sin \omega_r t$$

Step 3

$$\eta_r(0) = ? \text{ and } \dot{\eta}_r(0) = ?$$

Initial conditions $u(x,0) = \sum_{r=1}^{\infty} \phi_r(x) \eta_r(0) = Ax(1 - \frac{x}{L})$

$$\sum_{r=1}^{\infty} \int_0^L \phi_s(x) m(x) \phi_r(x) \eta_r(0) dx = \int_0^L m(x) \phi_s(x) Ax(1 - \frac{x}{L}) dx$$

$$\eta_r(0) = \int_0^L \frac{m \sqrt{2}}{\sqrt{mL}} \sin \frac{\lambda r x}{L} Ax(1 - \frac{x}{L}) dx = \frac{-2L^2}{(\lambda r)^3} [(-1)^r - 1] \frac{m \sqrt{2}}{\sqrt{mL}} A$$

$$= 0 \text{ for } r=2,4,\dots$$

$$= \sqrt{\frac{2}{mL}} \cdot \frac{4mL^3 A}{\lambda r^3} \quad r=1,3,\dots$$

and $\dot{\eta}_r(0) = 0$ as $\dot{u}(x,0) = 0$

$$\therefore u(x,t) = \sum_r \phi_r(x) \eta_r(t) = \sum_r \frac{2}{mL} \cdot \frac{4mL^3 A}{\lambda r^3} \cos \frac{\lambda r}{L} \sqrt{\frac{EA}{m}} \frac{\sin \lambda r x}{L}$$

$$= \sum_{r=1,3,\dots} \frac{8L^2 A}{\lambda^3 r^3} \cos \omega_r t \sin \frac{\lambda r x}{L}$$