

n. of motion of an Euler-Bernoulli beam:

$$\frac{\partial^2}{\partial x^2} \left[EI(x) \frac{\partial^2 y}{\partial x^2} \right] + m(x) \frac{\partial^3 y}{\partial x^3} = f(x, t) \quad (1)$$

Modal expansion theorem $y(x, t) = \sum_{v=1}^{\infty} \varphi_v(x) \eta_v(t) \quad (2)$

Substituting eqn. (2) into eqn. (1), we get,

$$\frac{d^2}{dx^2} \left[EI(x) \sum \varphi_v''(x) \eta_v(t) \right] + m(x) \sum \varphi_v(x) \ddot{\eta}_v(t) = f(x, t)$$

Multiplying both sides by $\varphi_s(x)$ and integrating, we get

$$\int_0^L \varphi_s(x) \frac{d^2}{dx^2} \left[EI(x) \sum \varphi_v''(x) \eta_v(t) \right] dx + \int_0^L m(x) \varphi_s(x) \sum \varphi_v(x) \ddot{\eta}_v(t) dx \\ = \int_0^L \varphi_s(x) f(x, t) dx$$

$$\varphi_s(x) \frac{d}{dx} \left[EI(x) \sum \varphi_v''(x) \eta_v(t) \right] \Big|_0^L - \int_0^L \frac{d\varphi_s}{dx} \cdot \frac{d}{dx} \left[EI(x) \frac{d^2 \varphi_v}{dx^2} \eta_v(t) \right] dx \\ + \ddot{\eta}_v(t) = \int_0^L \varphi_s(x) f(x, t) dx$$

$$\Rightarrow - \frac{d\varphi_s}{dx} \cdot \left[EI(x) \frac{d^2 \varphi_v}{dx^2} \right] \Big|_0^L + \int_0^L \frac{d^2 \varphi_s}{dx^2} \left[EI(x) \frac{d^2 \varphi_v}{dx^2} \eta_v(t) \right] dx + \ddot{\eta}_v(t) = \int_0^L \varphi_s(x) f(x, t) dx \\ = 0 \quad (\text{for all BCs}) \quad = \omega_v^2 \eta_v(t) \quad (\text{for } s=v) \\ = 0 \quad (\text{for all other } v \neq s)$$

$$\Rightarrow \int_0^L \omega_v^2 \eta_v + \ddot{\eta}_v = \int_0^L \varphi_v(x) f(x, t) dx = N_v(t)$$

Similar to 1D SDOF eqn 2.
can be solved accordingly.



A simply-supported beam is subjected to initial displacement given as,
 $y(x,0) = A \left(\frac{x}{L} - \frac{2x^3}{L^3} + \frac{x^4}{L^4} \right)$

at expansion theorem:

$$y(x,t) = \sum \varphi_r(x) \eta_r(t)$$

$$y(x,0) = \sum \varphi_r(x) \eta_r(0)$$

Multiplying both sides of the above expression with $\varphi_s(x) dx$ and integrating, we get

$$\int_0^L \varphi_s(x) m y(x,0) dx = \sum_r \int_0^L \varphi_s(x) \varphi_r(x) m dx \eta_r(0)$$

$$\Rightarrow \eta_r(0) = \int_0^L \varphi_s(x) m y(x,0) dx \quad \begin{aligned} & \eta_r(0) \text{ for } r=s \\ & = 0 \text{ for } r \neq s. \end{aligned}$$

where, $\varphi_r(x)$ (for simply-supported Euler-Bernoulli beam) = $\sqrt{\frac{2}{mL}} \sin \frac{\pi rx}{L}$.

$$\therefore \eta_r(0) = \int_0^L m \sqrt{\frac{2}{mL}} \sin \frac{\pi rx}{L} \cdot A \left(\frac{x}{L} - \frac{2x^3}{L^3} + \frac{x^4}{L^4} \right) dx \quad (1)$$

Eqn. of motion of an Euler-Bernoulli beam $EI \frac{\partial^4 y}{\partial x^4} + f(x,t) m \frac{\partial^2 y}{\partial t^2}$

After modal analysis $i\dot{\eta}_r + \omega_r^2 \eta_r = 0$

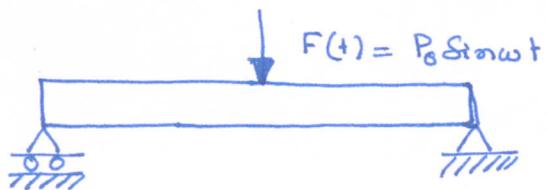
$$\Rightarrow \eta_r(t) = \eta_r(0) \cos \omega_r t + \frac{i\dot{\eta}_r(0)}{\omega_r} \sin \omega_r t$$

$$= \eta_r(0) \cos \left(\frac{\pi r}{L} \right)^2 \sqrt{\frac{EI}{m}} t$$

given by (1)

$$y(x,t) = \sum_{r=1}^{\infty} \sqrt{\frac{2}{mL}} \sin \frac{\pi rx}{L} \cdot \eta_r(0) \cos \left(\frac{\pi r}{L} \right)^2 \sqrt{\frac{EI}{m}} t \quad \underline{\underline{\text{Ans}}}$$

Example



Eqn. of motion

$$EI \frac{\partial^4 y}{\partial x^4} + m \frac{\partial^2 y}{\partial t^2} = f(x, t) = P_0 \delta(x - \frac{L}{2}) \sin \omega t$$

For a simply-supported Euler-Bernoulli beam:

$$\text{Mass normalized mode shape} = \sqrt{\frac{2}{mL}} \sin \frac{\pi n x}{L}$$

$$\text{Natural frequency } \omega_r = \left(\frac{\pi n}{L}\right)^2 \sqrt{\frac{EI}{m}}$$

After modal analysis:

$$\ddot{\eta}_r + \omega_r^2 \eta_r = \int_0^L \varphi_r f(x, t) dx$$

$$= P_0 \sqrt{\frac{2}{mL}} \int_0^L \sin \frac{\pi n x}{L} \cdot \delta(x - \frac{L}{2}) dx \sin \omega t$$

$$= P_0 \sqrt{\frac{2}{mL}} \sin \omega t \sin \frac{\pi n}{2} = A_0 \sin \omega t$$

$$\text{where } A_0 = P_0 \sqrt{\frac{2}{mL}} \sin \frac{\pi n}{2}$$

$$\therefore \eta_r(t) = A_0 H(\omega) \sin \omega t = \frac{A_0 \sin \omega t}{\omega_r^2 - \omega^2} = \frac{A_0}{\omega_r^2 - \omega^2} \sin \omega t$$

$$\therefore y(x, t) = \sqrt{\frac{2}{mL}} \cdot P_0 \cdot \underbrace{\sqrt{\frac{2}{mL}}}_{\text{cancel}} \sin \frac{\pi n}{2} \cdot \underbrace{\frac{1}{\omega_r^2 - \omega^2}}_{\text{cancel}} \sin \omega t$$

$$= \frac{2P_0}{mL} \cdot \frac{1}{\omega_r^2 - \omega^2} \sin \omega t \quad n = 1, 3, \dots$$

$$= - \frac{2P_0}{mL} \underbrace{\left[\frac{1}{\omega_r^2 - \omega^2} \right]}_{\text{cancel}} \sin \omega t \quad n = 2, 4, \dots$$