AE31002 Aerospace Structural Dynamics Laplace Transformation in Vibration Analysis

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Advantages

- It can handle discontinuous functions.
- It takes into account initial conditions automatically.

Laplace transformation of x(t), is written symbolically as $\bar{x}(s) = \mathcal{L}x(t)$ is defined by definite integral

$$\bar{x}(s) = \mathcal{L}x(t) = \int_0^\infty e^{-st} x(t) dt$$

where s is in general a complex quantity (s = x + iy) referred to as a subsidiary variable. The function e^{-st} is known as the kernel of the transformation.

For the SDOF system we need

$$\mathcal{L}\frac{dx(t)}{dt} = \int_0^\infty e^{-st} \frac{dx(t)}{dt} dt = s\bar{x}(s) - x(0)$$

and

$$\mathcal{L}\frac{d^2x(t)}{dt^2} = \int_0^\infty e^{-st} \frac{d^2x(t)}{dt^2} dt = s^2 \bar{x}(s) - sx(0) - \dot{x}(0)$$

Laplace transformation of the excitation function is

$$\bar{F}(s) = \mathcal{L}F(t) = \int_0^\infty e^{-st} F(t) dt$$
(1)

Applying to the damped SDOF system and rearranging

$$(ms^{2} + cs + k)\bar{x}(s) = \bar{F}(s) + m\dot{x}(0) + (ms + c)x(0)$$

Now ignoring the homogeneous solution i.e., letting $x(0) = \dot{x}(0) = 0$, ratio of transformed excitation to transformed response (generalised impedance of the system)

$$\bar{Z}(s) = rac{\bar{F}(s)}{\bar{x}(s)} = ms^2 + cs + k$$

Similarly the **transfer function** or the system function $\overline{G}(s)$ is expressed as

$$\bar{G}(s) = \frac{\bar{x}(s)}{\bar{F}(s)} = \frac{1}{ms^2 + cs + k} = \frac{1}{m(s^2 + 2\xi\omega_n s + \omega_n^2)}$$
(2)

The above equation may be written as

$$\bar{x}(s) = \bar{G}(s) \ \bar{F}(s)$$
 (3)

so that the **transfer function can be regarded as an algebraic operator that operates on the transformed excitation** to yield the transformed response.

Steps to follow are, to find out the equations as numbered $(1) \rightarrow (2) \rightarrow (3)$. After finding out the value of $\bar{x}(s)$ an inverse Laplace transformation is required to find out the response in real time t, i.e.,

$$x(t) = \mathcal{L}^{-1}\bar{x}(s) = \mathcal{L}^{-1}\bar{G}(s) \ \bar{F}(s)$$

Find out the response of a damped SDOF system under a unit impulse load.

Laplace transformation of unit impulse is

$$\bar{F}(s) = \mathcal{L}F(t) = 1$$

Laplace transformed response, $\bar{x}(s)$, in terms of transfer function $\bar{G}(s)$

$$\bar{x}(s) = \bar{G}(s) = \frac{1}{m(s^2 + 2\xi\omega_n s + \omega_n^2)}$$
$$\Rightarrow \bar{x}(s) = \frac{1}{2i\omega_d m} \left(\frac{1}{s + \xi\omega_n - i\omega_d} - \frac{1}{s + \xi\omega_n + i\omega_d}\right)$$

Response after the inverse Laplace transformation

$$\Rightarrow x(t) = \mathcal{L}^{-1}\bar{x}(s) = \frac{1}{2i\omega_d m} \left[e^{-(\xi\omega_n - i\omega_d)t} - e^{-(\xi\omega_n + i\omega_d)t} \right]$$
$$\Rightarrow x(t) = \frac{1}{m\omega_d} e^{-\xi\omega_n t} \sin \omega_d t$$

Find out the response of a damped SDOF system under a suddenly applied load of unit magnitude (step function with $F_0=1$).

$$\bar{F}(s) = \mathcal{L}F(t) = \int_0^\infty e^{-st} \cdot 1 \cdot dt = \frac{e^{-st}}{-s} \bigg|_0^\infty = \frac{1}{s}$$
$$\bar{x}(s) = \bar{G}(s)\bar{F}(s) = \frac{\bar{G}(s)}{s} = \frac{1}{ms(s^2 + 2\xi\omega_n s + \omega_n^2)}$$

Now following a standard process of inverse Laplace transformation, i.e., by factorising the denominator we get

$$\bar{x}(s) = \frac{1}{m\omega_n^2} \left[\frac{1}{s} - \frac{\xi\omega_n + i\omega_d}{2i\omega_d} \frac{1}{s + \xi\omega_n - i\omega_d} + \frac{\xi\omega_n - i\omega_d}{2i\omega_d} \frac{1}{s + \xi\omega_n + i\omega_d} \right]$$

Since

$$\mathcal{L}^{-1}\frac{1}{s-\alpha}=e^{\alpha t}$$

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$$\mathcal{L}^{-1}\bar{x}(s) = \frac{1}{k} \left[1 - \frac{\xi \omega_n + i\omega_d}{2i\omega_d} e^{-(\xi \omega_n - i\omega_d)t} + \frac{\xi \omega_n - i\omega_d}{2i\omega_d} e^{-(\xi \omega_n + i\omega_d)t} \right]$$

The above equation can be simplified to

$$x(t) = \mathcal{L}^{-1} ar{x}(s) = rac{1}{k} \left[1 - rac{1}{\sqrt{(1-\xi^2)}} e^{-\xi \omega_n t} cos \left(\omega_d t - \psi
ight)
ight]$$

where $tan \ \psi = rac{\xi}{\sqrt{(1-\xi^2)}}$

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Obtain the response of the system $m\ddot{x} + kx = kf(t)$ by the Laplace transform method. Assume no damping and consider f(t) = 0 for t<0.



$$\begin{aligned} x(t) &= \mathcal{L}^{-1}\bar{G}(s)\bar{F}(s) \\ x(t) &= \frac{kA}{T}\mathcal{L}^{-1}\left[\frac{1}{ms^2 + k} \cdot \frac{1}{s^2}\right] = \frac{\omega_n^2 A}{T}\mathcal{L}^{-1}\left[\frac{1}{s^2 + \omega_n^2} \cdot \frac{1}{s^2}\right] \\ x(t) &= = \frac{A}{T}\mathcal{L}^{-1}\left[\frac{1}{s^2} - \frac{1}{s^2 + \omega_n^2}\right] \\ x(t) &= = \frac{A}{T\omega_n}\left(\omega_n t - \sin\omega_n t\right), \quad t > 0 \end{aligned}$$

Applying to the damped SDOF system and rearranging

$$(ms^{2} + cs + k)\bar{x}(s) = \bar{F}(s) + m\dot{x}(0) + (ms + c)x(0)$$

Considering the initial boundary condition as $x(0) = x_0$ and $\dot{x}(0) = v_0$ we obtain the transformed response in the form

$$\bar{x}(s) = \frac{\bar{F}(s)}{m(s^2 + 2\xi\omega_n + \omega_n^2)} + \frac{s + 2\xi\omega_n}{s^2 + 2\xi\omega_n + \omega_n^2} x_0 + \frac{1}{s^2 + 2\xi\omega_n + \omega_n^2} v_0$$

It can be shown by inverse Laplace transform that

$$\begin{aligned} \mathbf{x}(t) &= \frac{1}{m\omega_d} \int_0^t F(\tau) e^{-\xi\omega_n(t-\tau)} \sin \omega_d(t-\tau) d\tau \\ &+ \frac{x_0}{\sqrt{1-\xi^2}} e^{-\xi\omega_n t} \cos(\omega_d t - \psi) + \frac{v_0}{\omega_d} e^{-\xi\omega_n t} \sin \omega_d t \end{aligned}$$